

New Superconformal Chern-Simons Theories

Imperial College July 2008

Kazuo Hosomichi, KML, Sangmin Lee,
Sungjay Lee, Jaemo Park

M2, D3, M5 Branes

Maximally Supersymmetric Conformal Field Theory

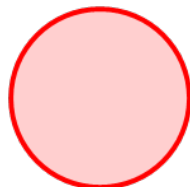
M2	2+1 dim	SO(8)	AdS ₄ x S ⁷	Lie group	$N^{\frac{3}{2}}$
D3	3+1 dim	SO(6)	AdS ₅ x S ⁵	Lie group	N^2
M5	5+1 dim (2,0)	SO(5)	AdS ₇ x S ⁴	ADE group	N^3

Strong Coupling Limit of Type IIA Theory = M Theory

D2 Brane \rightarrow M2 Brane

$$\Phi_I, I = 1, 2, 3 \dots 7$$

$$\mathcal{L} = -\frac{1}{4g^2} \text{tr} \left(F_{\mu\nu} F^{\mu\nu} - 2D_\mu \Phi_I D^\mu \Phi_I + \dots \right)$$



$$\text{length} \sim \frac{1}{g^2}$$

$$g^2 \sim \frac{R_{11}^2}{\ell_P^3}$$

Superconformal Chern-Simons Theories as Theories on M2 Branes

- Bagger-Lambert, Gustavsson $\mathcal{N} = 8$
SU(2)xSU(2), SO(8)R symmetry

- Aharony-Bergman-Jafferis-Maldacena
U(N)xU(N), Chern-Simons level k

$$\mathcal{N} = 6$$

$$\mathbb{C}^4 / \mathbb{Z}_k$$

$$k = 1, 2$$

N=3 Superconformal Theories

- Arbitrary Gauge Group
 - Arbitrary Representation
 - Parity Broken
 - U(1)xU(1): N=4 Theory
- Kapustin and Strassler (1999)

Spin

• 1	:	1
• $\frac{1}{2}$:	3
• 0	:	3
• $-\frac{1}{2}$:	1
• -1	:	0

$$\mathcal{L} \sim k A dA + (d\phi)^2 + \psi d\psi + \frac{1}{k} \phi^2 \psi^2 + \frac{1}{k^2} \phi^6$$

$N=3$ to $N=4,5,6,8$

- Gaiotto-Witten: (HLLP)
 $N=1 + SO(3) \text{ global} \Rightarrow SO(4) \text{ R}$
 $+ \text{twisted hyper} + \dots \Rightarrow 4,5,6,8$
- ABJM: (Schnabl, Tachikawa)
 $N=3 \Rightarrow 4,5,6,8$
- Deformed 3-Algebra: BL $\Rightarrow N=6$

3-dim Gaiotto-Witten Theory

- Start with N=1 Theory

- Matter: hypermultiplet

$$(q_{\alpha}^A, \psi_{\dot{\alpha}}^A)$$

- SO(3) global symmetry:
matter in 2-dim representation

- Gauge Groups: a subgroup of Sp(2n)

$$\omega_{AB} \quad (t^m)^A_B$$

Require $SO(3) \Rightarrow SU(2) \times SU(2) \times R$

- Write the most general N=1 superconformal theory with $SO(3)$ symmetry

$$((A_m)_\mu, \chi^m) \quad (q_\alpha^A, \psi_{\dot{\alpha}}^A, F_\alpha^A)$$

- Chern-Simons Term

$$\mathcal{L}_{\text{CS}} = \frac{\varepsilon^{\mu\nu\lambda}}{4\pi} k_{mn} A_\mu^m \partial_\nu A_\lambda^n + \frac{\varepsilon^{\mu\nu\lambda}}{12\pi} f_{mnp} A_\mu^m A_\nu^n A_\lambda^p + \frac{ik_{mn}}{4\pi} \chi^m \chi^n,$$

- Matter Lagrangian

$$\begin{aligned} \mathcal{L}_{\text{kin}} &= \frac{1}{2} \left(-D\bar{q}_A^\alpha Dq_\alpha^A + \bar{F}_A^\alpha F_\alpha^A + i\bar{\psi}_A^{\dot{\alpha}} \not{D} \psi_{\dot{\alpha}}^A - i\bar{\psi}_A^{\dot{\alpha}} \chi_{\dot{B}}^A q_\alpha^B + i\bar{q}_A^\alpha \chi_{\dot{B}}^A \psi_{\dot{\alpha}}^B \right) \\ &= \frac{1}{2} \omega_{AB} \left[\epsilon^{\alpha\beta} \left(-Dq_\alpha^A Dq_\beta^B + F_\alpha^A F_\beta^B \right) + i\epsilon^{\dot{\alpha}\dot{\beta}} \psi_{\dot{\alpha}}^A \not{D} \psi_{\dot{\beta}}^B \right] - i\epsilon^{\dot{\alpha}\dot{\beta}} \psi_{\dot{\alpha}}^A \chi_{\dot{B}}^A q_{\dot{\beta}}^B, \end{aligned}$$

- Superpotential W

$$W = \frac{\pi}{4} T_{AB,CD} \epsilon^{\alpha\beta} \epsilon^{\gamma\delta} q_{\alpha}^A q_{\beta}^B q_{\gamma}^C q_{\delta}^D.$$

- W- Lagrangian

$$-\pi T_{AB,CD} \left(\epsilon^{\alpha\beta} \epsilon^{\gamma\delta} F_{\alpha}^A q_{\beta}^B q_{\gamma}^C q_{\delta}^D + \frac{i}{2} \epsilon^{\dot{\alpha}\dot{\beta}} \epsilon^{\gamma\delta} \psi_{\dot{\alpha}}^A \dot{\psi}_{\dot{\beta}}^B q_{\gamma}^C q_{\delta}^D + i \epsilon^{\dot{\alpha}\beta} \epsilon^{\dot{\gamma}\delta} \psi_{\dot{\alpha}}^A q_{\beta}^B \dot{\psi}_{\dot{\gamma}}^C q_{\delta}^D \right)$$

- Integrate over gaugino & SO(4)-violating terms is

$$-i\pi q_{\alpha}^A q_{\beta}^B \dot{\psi}_{\dot{\gamma}}^C \dot{\psi}_{\dot{\delta}}^D \left(k_{mn} t_{AC}^m t_{BD}^n \epsilon^{\alpha\dot{\gamma}} \epsilon^{\beta\dot{\delta}} + \frac{1}{2} T_{AB,CD} \epsilon^{\alpha\beta} \epsilon^{\dot{\gamma}\dot{\delta}} - T_{AC,BD} \epsilon^{\alpha\dot{\gamma}} \epsilon^{\beta\dot{\delta}} \right)$$

- SO(4) R-Symmetry

$$T_{AC,BD} + T_{BC,AD} + k_{mn} (t_{AC}^m t_{BD}^n + t_{BC}^m t_{AD}^n) = 0.$$

- Fundamental Identity:
constraints on gauge generators

$$k_{mn} t^m_{(AB} t^n_{C)D} = 0,$$

- Coupling constants:

$$T_{AB,CD} = \frac{1}{3} k_{mn} (t^m_{AC} t^n_{BD} - t^m_{BC} t^n_{AD})$$

- SU(2)xSU(2) R-symmetry:

$$\delta\phi_{\alpha}^A = i\eta\psi_{\dot{\alpha}}^A \rightarrow \delta\phi_{\alpha}^A = i\eta_{\alpha}^{\dot{\alpha}}\psi_{\dot{\alpha}}^A$$

N=4 Lagrangian and Supersymmetric Transformation

$$\begin{aligned}\mathcal{L} = & \frac{\varepsilon^{\mu\nu\lambda}}{4\pi} \left(k_{mn} A_\mu^m \partial_\nu A_\lambda^n + \frac{1}{3} f_{mnp} A_\mu^m A_\nu^n A_\lambda^p \right) \\ & + \frac{1}{2} \omega_{AB} \left(-\epsilon^{\alpha\beta} D q_\alpha^A D q_\beta^B + i \epsilon^{\dot{\alpha}\dot{\beta}} \psi_{\dot{\alpha}}^A \not{D} \psi_{\dot{\beta}}^B \right) \\ & - i\pi k_{mn} \epsilon^{\alpha\beta} \epsilon^{\dot{\gamma}\dot{\delta}} j_{\alpha\dot{\gamma}}^m j_{\beta\dot{\delta}}^n - \frac{\pi^2}{6} f_{mnp} (\mu^m)^\alpha{}_\beta (\mu^n)^\beta{}_\gamma (\mu^p)^\gamma{}_\alpha\end{aligned}$$

$$\mu_{\alpha\beta}^m = t_{AB}^m q_\alpha^A q_\beta^B$$

$$\begin{aligned}\delta q_\alpha^A &= i\eta_\alpha{}^{\dot{\alpha}} \psi_{\dot{\alpha}}^A, & \delta A_\mu^m &= 2\pi i \eta^{\alpha\dot{\alpha}} \gamma_\mu j_{\alpha\dot{\alpha}}^m, \\ \delta \psi_{\dot{\alpha}}^A &= \left[\not{D} q_\alpha^A + \frac{2\pi}{3} (t_m)^A{}_B q_\beta^B (\mu^m)^\beta{}_\alpha \right] \eta^{\alpha}{}_{\dot{\alpha}}.\end{aligned}$$

$$j_{\alpha\dot{\gamma}}^m \equiv q_\alpha^A t_{AC}^m \psi_{\dot{\gamma}}^C$$

Lie Super-Algebra

$$[M^m, M^n] = f^{mn}_p M^p$$

$$[M^m, Q_A] = Q_B (t^m)^B_A$$

$$\{Q_A, Q_B\} = t^m_{AB} M_m$$

Jacobi Identity -> Fundamental Identity

Gauge Group and Matter Representation

Classification of Super Lie Algebra

- $U(N) \times U(M)$
 - $O(N) \times Sp(2M)$
 - $SO(7) \times Sp(2)$ 8-dim spinor
 - $G_2 \times Sp(2)$
 - $SU(2) \times SU(2) \times SU(2)$
 - $SU(N)$, adjoint
-
- The Chern-Simons level $k, -k$

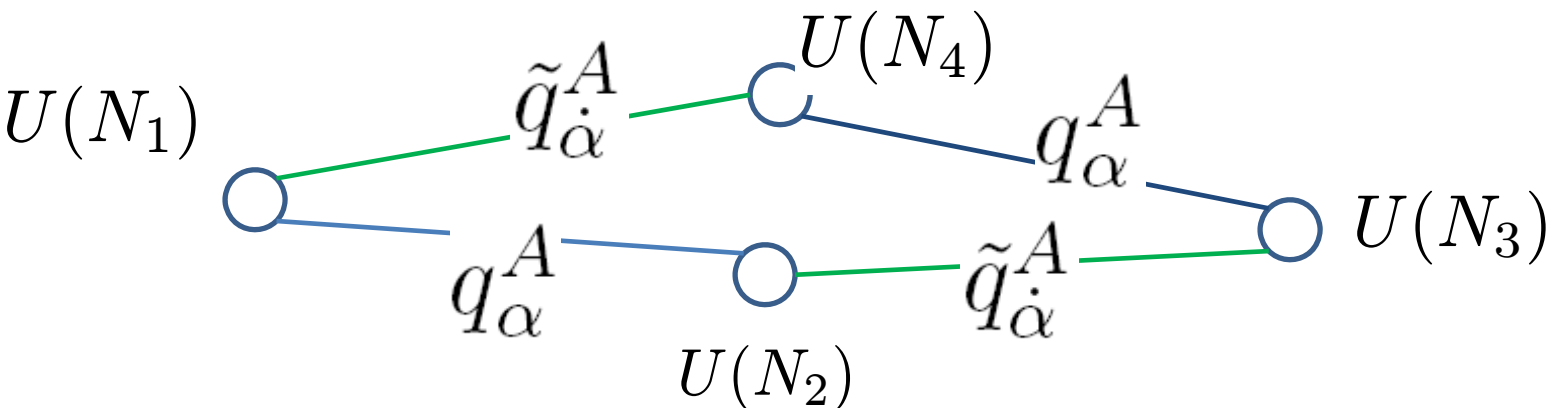
• Our Generalization

- Twisted Hyper-Multiplet

$$\tilde{q}_{\dot{\alpha}}^A, \tilde{\psi}_{\alpha}^A$$

$$\tilde{\mu}_{\dot{\alpha}\dot{\beta}}^m \equiv \tilde{t}_{AB}^m \tilde{q}_{\dot{\alpha}}^A \tilde{q}_{\dot{\beta}}^B \quad \tilde{J}_{\dot{\alpha}\alpha}^m \equiv \tilde{q}_{\dot{\alpha}}^A \tilde{t}_{AB}^m \tilde{\psi}_{\alpha}^B$$

- Linear Quiver Type Theory



$$\begin{aligned}
\mathcal{L} = & \frac{\varepsilon^{\mu\nu\lambda}}{4\pi} \left(k_{mn} A_\mu^m \partial_\nu A_\lambda^n + \frac{1}{3} f_{mnp} A_\mu^m A_\nu^n A_\lambda^p \right) \\
& + \frac{1}{2} \omega_{AB} \left(-\epsilon^{\alpha\beta} D q_\alpha^A D q_\beta^B + i \epsilon^{\dot{\alpha}\dot{\beta}} \psi_{\dot{\alpha}}^A \not{D} \psi_{\dot{\beta}}^B \right) + \frac{1}{2} \omega_{AB} \left(-\epsilon^{\dot{\alpha}\dot{\beta}} D \tilde{q}_{\dot{\alpha}}^A D \tilde{q}_{\dot{\beta}}^B + i \epsilon^{\alpha\beta} \tilde{\psi}_\alpha^A \not{D} \tilde{\psi}_\beta^B \right) \\
& - i \pi k_{mn} \epsilon^{\alpha\beta} \epsilon^{\dot{\gamma}\dot{\delta}} j_{\alpha\dot{\gamma}}^m j_{\beta\dot{\delta}}^n - i \pi k_{mn} \epsilon^{\dot{\alpha}\dot{\beta}} \epsilon^{\gamma\delta} \tilde{j}_{\dot{\alpha}\gamma}^m \tilde{j}_{\dot{\beta}\delta}^n + 4 \pi i k_{mn} \epsilon^{\alpha\gamma} \epsilon^{\dot{\beta}\dot{\delta}} j_{\alpha\dot{\beta}}^m \tilde{j}_{\dot{\delta}\gamma}^n \\
& + i \pi k_{mn} \left(\epsilon^{\dot{\alpha}\dot{\gamma}} \epsilon^{\beta\dot{\delta}} \tilde{\mu}_{\dot{\alpha}\dot{\beta}}^m \psi_{\dot{\gamma}}^A t_{AB}^n \psi_{\dot{\delta}}^B + \epsilon^{\alpha\gamma} \epsilon^{\beta\delta} \mu_{\alpha\beta}^m \tilde{\psi}_\gamma^A \tilde{t}_{AB}^n \tilde{\psi}_\delta^B \right) \\
& - \frac{\pi^2}{6} f_{mnp} (\mu^m)^\alpha{}_\beta (\mu^n)^\beta{}_\gamma (\mu^p)^\gamma{}_\alpha - \frac{\pi^2}{6} f_{mnp} (\tilde{\mu}^m)^{\dot{\alpha}}{}_{\dot{\beta}} (\tilde{\mu}^n)^{\dot{\beta}}{}_{\dot{\gamma}} (\tilde{\mu}^p)^{\dot{\gamma}}{}_{\dot{\alpha}} \\
& + \pi^2 (\tilde{\mu}^{mn})^{\dot{\gamma}}{}_{\dot{\gamma}} (\mu_m)^\alpha{}_\beta (\mu_n)^\beta{}_\alpha + \pi^2 (\mu^{mn})^\gamma{}_\gamma (\tilde{\mu}_m)^{\dot{\alpha}}{}_{\dot{\beta}} (\tilde{\mu}_n)^{\dot{\beta}}{}_{\dot{\alpha}}.
\end{aligned} \tag{2.4}$$

$$\begin{aligned}
\delta q_\alpha^A &= +i \eta_\alpha^{\dot{\alpha}} \psi_{\dot{\alpha}}^A, \quad \delta \tilde{q}_{\dot{\alpha}}^A = -i \eta_{\dot{\alpha}}^\alpha \tilde{\psi}_\alpha^A, \quad \delta A_\mu^m = 2\pi i \eta^{\alpha\dot{\alpha}} \gamma_\mu (j_{\alpha\dot{\alpha}}^m - \tilde{j}_{\dot{\alpha}\alpha}^m), \\
\delta \psi_{\dot{\alpha}}^A &= + \left[\not{D} q_\alpha^A + \frac{2\pi}{3} (t_m)^A{}_B q_\beta^B (\mu^m)^\beta{}_\alpha \right] \eta^\alpha{}_{\dot{\alpha}} - 2\pi (t_m)^A{}_B q_\beta^B (\tilde{\mu}^m)^{\dot{\beta}}{}_{\dot{\alpha}} \eta^\beta{}_{\dot{\beta}}, \\
\delta \tilde{\psi}_\alpha^A &= - \left[\not{D} \tilde{q}_{\dot{\alpha}}^A + \frac{2\pi}{3} (\tilde{t}_m)^A{}_B \tilde{q}_{\dot{\beta}}^B (\tilde{\mu}^m)^{\dot{\beta}}{}_{\dot{\alpha}} \right] \eta_\alpha^{\dot{\alpha}} + 2\pi (\tilde{t}_m)^A{}_B \tilde{q}_{\dot{\beta}}^B (\mu^m)^\beta{}_\alpha \eta_\beta^{\dot{\beta}}.
\end{aligned}$$

SU(2)xSU(2) gauge group: BLG Theory of N=8

Mass deformation

$$-\omega_{AB} \left(\frac{m^2}{2} \epsilon^{\alpha\beta} q_\alpha^A q_\beta^B + \frac{m^2}{2} \epsilon^{\dot{\alpha}\dot{\beta}} \tilde{q}_{\dot{\alpha}}^A \tilde{q}_{\dot{\beta}}^B + \frac{i}{2} m \epsilon^{\dot{\alpha}\dot{\beta}} \psi_{\dot{\alpha}}^A \psi_{\dot{\beta}}^B - \frac{i}{2} m \epsilon^{\alpha\beta} \tilde{\psi}_\alpha^A \tilde{\psi}_\beta^B \right) \\ - \frac{2\pi}{3} m k_{mn} \left((\mu^m)_{\alpha\beta} (\mu^n)^{\beta\alpha} - (\tilde{\mu}^m)_{\dot{\alpha}\dot{\beta}} (\tilde{\mu}^n)^{\dot{\beta}\dot{\alpha}} \right),$$

$$\delta_{\text{mass}} \psi_{\dot{\alpha}}^A = m q_\alpha^A \eta_{\dot{\alpha}}^\alpha, \quad \delta_{\text{mass}} \tilde{\psi}_\alpha^A = m \tilde{q}_{\dot{\alpha}}^A \eta_\alpha^{\dot{\alpha}}.$$

N=5

SCFT

$$t_{AB}^m = \tilde{t}_{AB}^m$$

$$G_1 \times G_2$$

Two Node System

$$\Phi_{\alpha}^A = \begin{pmatrix} q_{\alpha}^A \\ \tilde{q}_{\dot{\alpha}}^A \end{pmatrix}, \quad \Psi_{\alpha}^A = \begin{pmatrix} \tilde{\psi}_{\alpha}^A \\ \psi_{\dot{\alpha}}^A \end{pmatrix}$$

SU(2)xSU(2) => Sp(4) R-symmetry

$$C^{\alpha\beta} = \begin{pmatrix} \epsilon^{\alpha\beta} & 0 \\ 0 & \epsilon^{\dot{\alpha}\dot{\beta}} \end{pmatrix} \quad \mathcal{M}_{\alpha\beta}^m \equiv t_{AB}^m \Phi_{\alpha}^A \Phi_{\beta}^B \quad \mathcal{J}_{\alpha\beta}^m \equiv t_{AB}^m \Phi_{\alpha}^A \Psi_{\beta}^B$$

$$\begin{aligned} \mathcal{L} = & \frac{\varepsilon^{\mu\nu\lambda}}{4\pi} \left(k_{mn} A_m^m \partial_{\nu} A_{\lambda}^n + \frac{1}{3} f_{mnp} A_{\mu}^m A_{\nu}^n A_{\lambda}^p \right) \\ & + \frac{1}{2} \omega_{AB} C^{\alpha\beta} \left(-D\Phi_{\alpha}^A D\Phi_{\beta}^B + i\Psi_{\alpha}^A \not{D}\Psi_{\beta}^B \right) - i\pi k_{mn} C^{\alpha\beta} C^{\gamma\delta} \left(\mathcal{J}_{\alpha\gamma}^m \mathcal{J}_{\beta\delta}^n - 2\mathcal{J}_{\alpha\gamma}^m \mathcal{J}_{\delta\beta}^n \right) \\ & + \frac{2\pi^2}{15} f_{mnp} (\mathcal{M}^m)^{\alpha}_{\beta} (\mathcal{M}^n)^{\beta}_{\gamma} (\mathcal{M}^p)^{\gamma}_{\alpha} + \frac{3\pi^2}{5} (\mathcal{M}^{mn})^{\gamma}_{\gamma} (\mathcal{M}_m)^{\alpha}_{\beta} (\mathcal{M}_n)^{\beta}_{\alpha}, \end{aligned} \quad (2.13)$$

$$\begin{aligned} \delta\Phi_{\alpha}^A &= i\eta_{\alpha}^{\beta} \Psi_{\beta}^A, \quad \delta A_{\mu}^m = 2\pi i \eta^{\alpha\beta} \gamma_{\mu} \mathcal{J}_{\alpha\beta}^m, \\ \delta\Psi_{\alpha}^A &= \left[\not{D}\Phi_{\gamma}^A + \frac{2\pi}{3} (t_m)^A_B \Phi_{\beta}^B (\mathcal{M}^m)^{\beta}_{\gamma} \right] \eta^{\gamma}_{\alpha} - \frac{4\pi}{3} (t_m)^A_B \Phi_{\beta}^B (\mathcal{M}^m)^{\gamma}_{\alpha} \eta^{\beta}_{\gamma}. \end{aligned}$$

Sp(2M)x O(N), U(M)xU(N)

N=6 SCFT

Complex Representation

$$(t^A{}_B)_{\mathcal{N}=5} = \begin{pmatrix} t^A{}_B & 0 \\ 0 & -t_A{}^B \end{pmatrix}$$

$$(\Phi^A_\alpha)_{\mathcal{N}=5} = \begin{pmatrix} \Phi^A_\alpha \\ C_{\alpha\beta} \bar{\Phi}^\beta_A \end{pmatrix}, \quad (\Psi^A_\alpha)_{\mathcal{N}=5} = \begin{pmatrix} C_{\alpha\beta} \Psi^{\beta A} \\ -\bar{\Psi}_{\alpha A} \end{pmatrix}$$

$$(M^m)^\alpha{}_\beta \equiv \bar{\Phi}^\alpha_A (t^m)^A{}_B \Phi^B_\beta, \quad (M^{mn})^\alpha{}_\beta \equiv \bar{\Phi}^\alpha_A (t^m t^n)^A{}_B \Phi^B_\beta \\ (J^m)_{\alpha\beta} \equiv \Phi^A_\alpha (t^m)^A{}_B \bar{\Psi}_{\beta B}, \quad (\bar{J}^m)^{\alpha\beta} \equiv \bar{\Phi}^\alpha_A (t^m)^A{}_B \Psi^{\beta B}.$$

Sp(4) => SU(4) R-Symmetry

$$\begin{aligned}
\mathcal{L} = & \frac{\varepsilon^{\mu\nu\lambda}}{4\pi} \left(k_{mn} A_m^m \partial_\nu A_\lambda^n + \frac{1}{3} f_{mnp} A_\mu^m A_\nu^n A_\lambda^p \right) - D\bar{\Phi}_A^\alpha D\Phi_\alpha^A + i\bar{\Psi}_{\alpha A} \not{D}\Psi^{\alpha A} \\
& + i\pi \left[2(\bar{J}_m)^{\alpha\beta} (J_m)_{\alpha\beta} - 4(\bar{J}_m)^{\alpha\beta} (J^m)_{\beta\alpha} + \epsilon^{\alpha\beta\gamma\delta} (J_m)_{\alpha\beta} (J^m)_{\gamma\delta} + \epsilon_{\alpha\beta\gamma\delta} (\bar{J}_m)^{\alpha\beta} (\bar{J}^m)^{\gamma\delta} \right] \\
& - \frac{4\pi^2}{3} f_{mnp} (M^m)^\alpha_\beta (M^n)^\beta_\gamma (M^p)^\gamma_\alpha + 4\pi^2 (M^{mn})^\alpha_\beta (M_n)^\beta_\gamma (M_m)^\gamma_\alpha.
\end{aligned} \tag{3.13}$$

$$\begin{aligned}
\delta\Phi_\alpha^A &= -i\eta_{\alpha\beta} \Psi^{A\beta}, \quad \delta A_\mu^m = 2\pi i \left(\eta^{\alpha\beta} \gamma_\mu (J^m)_{\alpha\beta} + \eta_{\alpha\beta} \gamma_\mu (\bar{J}^m)^{\alpha\beta} \right), \\
\delta\Psi^{A\alpha} &= \left[\not{D}\Phi_\gamma^A - \frac{2\pi}{3} (t_m)^A_B \Phi_\beta^B (M^m)^\beta_\gamma \right] \eta^{\gamma\alpha} + \frac{4\pi}{3} (t_m)^A_B \Phi_\beta^B (M^m)^\alpha_\gamma \eta^{\gamma\beta} \\
&+ \frac{2\pi}{3} \epsilon^{\alpha\beta\gamma\delta} (t_m)^A_B \Phi_\beta^B (M^m)^\rho_\gamma \eta_{\delta\rho}.
\end{aligned}$$

U(M)xU(N): ABJM Case

U(1)xSp(2N): Something New

$$U(2N) \times U(2N) \rightarrow Sp(2N) \times O(2N)$$

M2 Branes on Orbifolds

$$\mathbb{C}^4 / \mathbb{Z}_k \quad (\phi_1, \phi_2, \phi_3, \phi_4)$$

$$\alpha : \phi_i \longrightarrow e^{2\pi i/k} \phi_i.$$

$$\mathbb{C}^4 / \hat{D}_{k+2}$$

$$\beta : (\phi_1, \phi_2, \phi_3, \phi_4) \longrightarrow (\phi_2^*, -\phi_1^*, \phi_4^*, -\phi_3^*),$$

$$\alpha^{2k} = 1, \quad \beta^2 = \alpha^k, \quad \beta \alpha \beta^{-1} = \alpha^{-1}.$$

Mass Deformation: supersymmetry preserved

- BLG: Supersymmetry Preserved,
 $SO(8) \text{ R} \rightarrow SO(4) \times SO(4)$
- Gaiotto-Witten: $SO(4) \text{ R}$ preserved
- $N=5$ Theories: $SO(5) \text{ R} \rightarrow SO(4)$
- $N=6$ Theories: $SO(6) \text{ R} \rightarrow SO(4) \times U(1)$

Conclusion

- New Superconformal Chern-Simons Theories are found.
- They deserve further exploration.
- There may be more such theories to be discovered.